Reverberation algorithms

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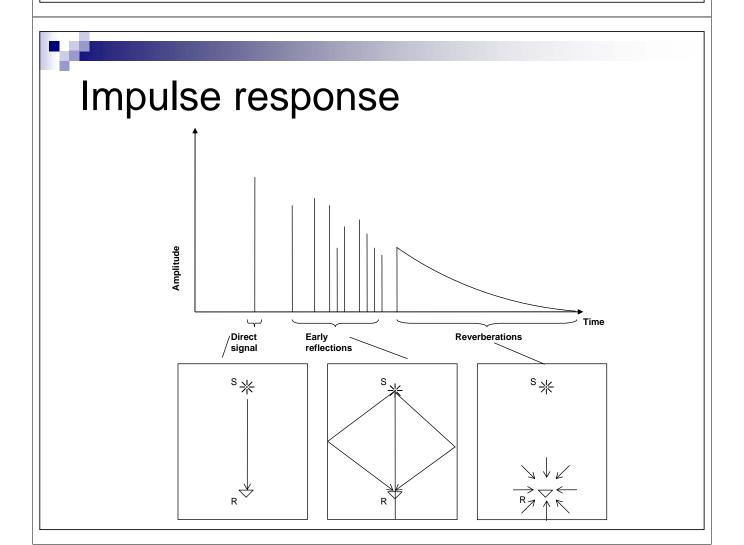
Summary

- The Reverb Problem
- Reverb Perception
- Acoustic impulse response:
 - □ Formation mechanisms
 - Parameters
- Early Reflections
- Late Reverb
- Numerical reverberation algorithms
 - □ Schroeder Reverbs
 - ☐ Feedback Delay Network (FDN) Reverberators
 - □ Waveguide Reverberators
- Geometrical reverberation algorithms



Impulse response

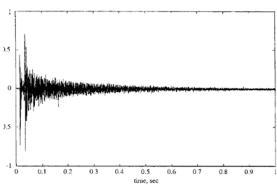
- The sounds we perceive heavily depend on the surrounding environment
- Environment-related sound changes are of convolutive origin (filtering)
 - □ Well-modeled by a space-varying impulse response

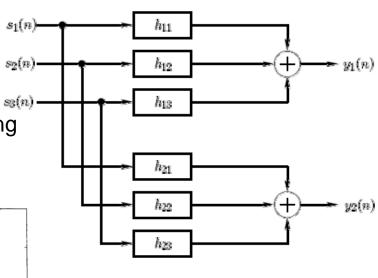


Reverberation tf Function

- Three sources, one listener (two ears)
- Filters should include pinnae filtering
- Filters change if anything in the room changes

(exact model)







Global descriptors

Energy decay curve (EDC)

$$\mathsf{EDC}(t) \stackrel{\Delta}{=} \int_t^\infty h^2(\tau) d\tau$$

- ☐ Introduced by Schroeder to define reverberation time
- □ It measures the total signal energy remaining in the reverberator's impulse response at time *t*
- ☐ It decays more smoothly than the impulse response, therefore it works better than the amplitude's envelope for defining the reverberation time
- ☐ In reverberant environments a large amount of the total energy is contained in the last portion of the impulse response
- Reverberation time

$$T60 = \{t : EDC(t) = EDC(0) - 60dB\}$$

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Global descriptors

The energy decay relief (EDR) generalizes the EDC to multiple frequency bands:

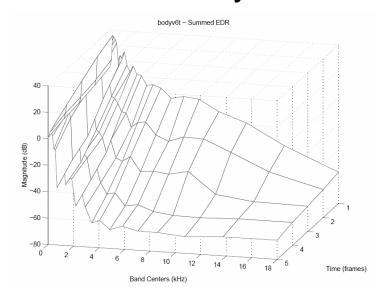
$$\mathsf{EDR}(t_n, f_k) \stackrel{\Delta}{=} \sum_{m=n}^{M} |H(m, k)|^2$$

where H(m,k) denotes bin k of the short-time Fourier transform (STFT) at time-frame m, and M is the number of frames.

- FFT window length $\approx 30 40$ ms
- EDR (t_n, f_k) = total signal energy remaining at time t_n sec in frequency band centered at f_k



EDR of a violin body



- Energy summed over frequency within each "critical band of hearing" (Bark band)
- Violin body = "small box reverberator"

Global descriptors

- In the room's transfer function we can single out resonant modes
- The spacing between two resonant modes is given by

$$\Delta f_{\rm max} \approx \frac{4}{T_r} \; {\rm Hz}.$$

which is valid above the threshold frequency

$$f_g \approx 2000 \sqrt{\frac{T_r}{V}} \text{ Hz.}$$

Global descriptors

Number of echoes in the impulse response before time t

$$N_t = \frac{4\pi (ct)^3}{3V}$$

■ Derivative of *N_t*:

$$\frac{dN_t}{dt} = \frac{4\pi c^3}{V}t^2$$

 Clarity index: ratio btw early reflections energy and late reverberation energy

$$C = 10 \log_{10} \left\{ \frac{\int_0^{80 \text{ ms}} p^2(t)dt}{\int_{80 \text{ ms}}^{\infty} p^2(t)dt} \right\} \text{ dB}$$

Implementation

Let $h_{ij}(n) = \text{impulse response from source } j$ to ear i. Then the output is given by $six\ convolutions$:

$$y_1(n) = (s_1 * h_{11})(n) + (s_2 * h_{12})(n) + (s_3 * h_{13})(n)$$

$$y_2(n) = (s_1 * h_{21})(n) + (s_2 * h_{22})(n) + (s_3 * h_{23})(n)$$

- ullet For small n, filters $h_{ij}(n)$ are sparse
- Tapped Delay Line (TDL) a natural choice

Transfer-function matrix:

$$\begin{bmatrix} Y_1(z) \\ Y_2(z) \end{bmatrix} = \begin{bmatrix} H_{11}(z) & H_{12}(z) & H_{13}(z) \\ H_{21}(z) & H_{22}(z) & H_{23}(z) \end{bmatrix} \begin{bmatrix} S_1(z) \\ S_2(z) \\ S_3(z) \end{bmatrix}$$



Complexity of Exact Reverberation

Example:

- Let $t_{60} = 1$ second
- $f_s = 50 \text{ kHz}$
- Each filter h_{ij} requires 50,000 multiplies and additions per sample, or 2.5 *billion* multiply-adds per second.
- Three sources and two listening points (ears) ⇒
 30 billion operations per second

Conclusion: Exact implementation of point-to-point transfer functions is too expensive for real-time computation.

Possibility of a Physical Reverb Model

In a complete physical model of a room,

- sources and listeners can be moved without affecting the room simulation itself,
- spatialized (in 3D) stereo output signals can be extracted using a "virtual dummy head"



How expensive is a room physical model?

- ullet Audio bandwidth $= 20 \text{ kHz} pprox 1/2 inch wavelength}$
- \bullet Spatial samples every 1/4 inch or less
- A 12'x12'x8' room requires > 100 million grid points
- ullet A lossless 3D finite difference model requires one multiply and 6 additions per grid point \Rightarrow 30 billion additions per second at $f_s=50~\mathrm{kHz}$
- A 100'x50'x20' concert hall requires more than 3 quadrillion operations per second

Conclusion: Fine-grained physical models are too expensive for real-time computation, especially for large halls.

Perceptual Aspects of Reverberation

Artificial reverberation is an unusually interesting signal processing problem:

- "Obvious" methods based on physical modeling or input-output modeling are too expensive
- We do not perceive the full complexity of reverberation
- What is important perceptually?
- How can we simulate only what is audible?



Perception of Echo Density and Mode Density

- For typical rooms
 - Echo density increases as t^2
 - Mode density increases as f^2
- Beyond some time, the echo density is so great that a stochastic process results
- Above some frequency, the mode density is so great that a random frequency response results
- There is no need to simulate many echoes per sample
- There is no need to implement more resonances than the ear can hear

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Proof that Echo Density Grows as Time Squared

Consider a single spherical wave produced from a point source in a rectangular room.

- Tesselate 3D space with copies of the original room
- Count rooms intersected by spherical wavefront



Proof that Mode Density Grows as Freq. Squared

The resonant modes of a rectangular room are given by

$$k^{2}(l,m,n) = k_{x}^{2}(l) + k_{y}^{2}(m) + k_{z}^{2}(n)$$

- $k_x(l) = l\pi/L_x = l$ th harmonic of the fundamental standing wave in the x
- ullet $L_x = \mbox{length of the room along } x$
- ullet Similarly for y and z
- \bullet Mode frequencies map to a uniform 3D Cartesian grid indexed by (l,m,n)
- Grid spacings are π/L_x , π/L_y , and π/L_z in x,y, and z, respectively.
- Spatial frequency k of mode (l,m,n)= distance from the (0,0,0) to (l,m,n)
- ullet Therefore, the number of room modes having a given spatial frequency grows as k^2

Early Reflections and Late Reverb

Based on limits of perception, the impulse response of a reverberant room can be divided into two segments

- Early reflections = relatively sparse first echoes
- Late reverberation—so densely populated with echoes that it is best to characterize the response statistically.



Early Reflections and Late Reverb

Similarly, the *frequency response* of a reverberant room can be divided into two segments.

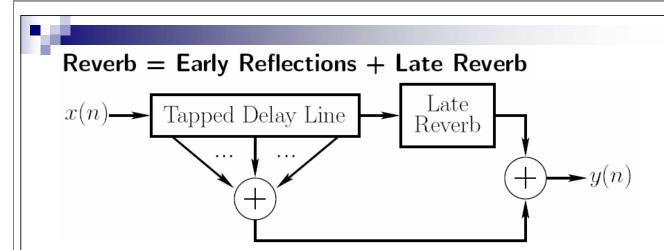
- Low-frequency sparse distribution of resonant modes
- Modes packed so densely that they merge to form a random frequency response with regular statistical properties

Perceptual Metrics for Ideal Reverberation

Some desirable controls for an artificial reverberator include

- ullet $t_{60}(f) = {\sf desired}$ reverberation time at each frequency
- $G^2(f) = \text{signal power gain at each frequency}$
- ullet C(f)= "clarity" = ratio of impulse-response energy in early reflections to that in the late reverb
- ullet ho(f)= inter-aural correlation coefficient at left and right ears

Perceptual studies indicate that reverberation time $t_{60}(f)$ should be independently adjustable in at least *three* frequency bands.



- TDL taps may include lowpass filters (air absorption, lossy reflections)
- Several taps may be fed to late reverb unit, especially if it takes a while to reach full density
- Some or all early reflections can usually be worked into the delay lines of the late-reverberation simulation (transposed tapped delay line)



Early Reflections

The "early reflections" portion of the impulse response is defined as everything up to the point at which a statistical description of the late reverb takes hold.

- Often taken to be the first 100ms
- Better to test for Gaussianness
 - Histogram test for sample amplitudes in 10ms windows
 - Exponential fit to EDC (Prony's method, matrix pencil method)
 - Crest factor test (peak/rms)



Early Reflections

- Typically implemented using tapped delay lines (TDL) (suggested by Schroeder in 1970 and implemented by Moorer in 1979)
- Early reflections should be spatialized (Kendall)
- Early reflections influence spatial impression



Late Reverberation

Desired Qualities:

- 1. a smooth (but not too smooth) decay, and
- 2. a smooth (but not too regular) frequency response.
- Exponential decay no problem
- Hard part is making it *smooth*
 - Must not have "flutter," "beating," or unnatural irregularities
 - Smooth decay generally results when the echo density is sufficiently high
 - Some short-term energy fluctuation is required for naturalness



Late Reverberation

Desired Qualities:

- A smooth *frequency response* has no large "gaps" or "hills"
 - Generally provided when the mode density is sufficiently large
 - Modes should be spread out uniformly
 - Modes may not be too regularly spaced, since audible periodicity in the time-domain can result

Late Reverberation

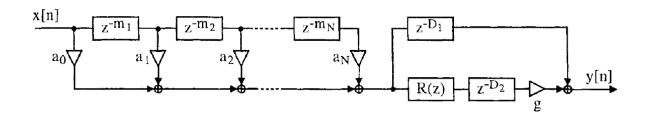
Desired Qualities:

- Moorer's ideal late reverb: exponentially decaying white noise
 - Good smoothness in both time and frequency domains
 - High frequencies need to decay faster than low frequencies
- Schroeder's rule of thumb for echo density in the late reverb is 1000 echoes per second or more
- For impulsive sounds, 10,000 echoes per second or more may be necessary for a smooth response

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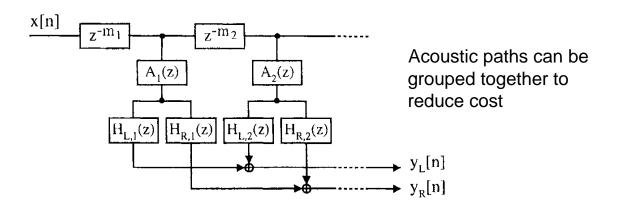
Moorer reverberator

 accounts for late reverberations by placing an IIR filter after the FIR filter (tapped delay line)



Binaural impulse response

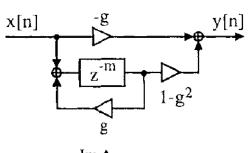
- Our sound perception is affected by our own body
 - ☐ Head Related Transfer Function (HRTF)

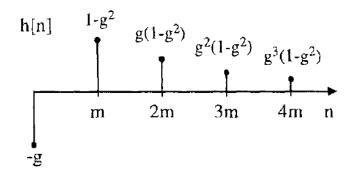


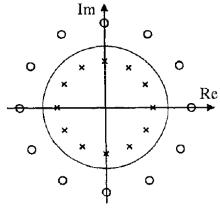
Comb filter h[n]y[n]x[n] g^2 g^3 3m 0 2m4m m $|H(e^{j\omega})|$ 1/(1-g)Re 1/(1+g)0 1/mT2/mT $\omega/2\pi$

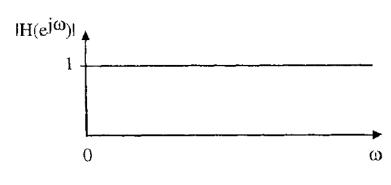
A.U. CIG

Allpass filter









Why Allpass?

- Allpass filters do not occur in natural reverberation!
- "Colorless reverberation" is an idealization only possible in the "virtual world"
- Perceptual factorization:

Coloration now orthogonal to decay time and echo density

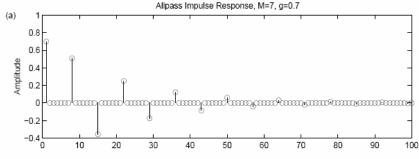
Are Allpasses Really Colorless?

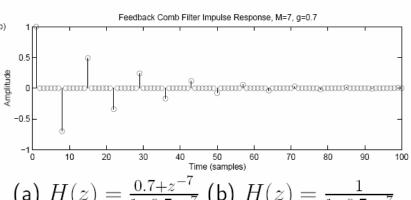
- Allpass impulse response only "colorless" when extremely short (less than 10 ms or so).
- Long allpass impulse responses sound like feedback comb-filters
- The difference between an allpass and feedback-comb-filter impulse response is *one echo*!



Are Allpasses Really Colorless?

Steady-state tones (sinusoids) really do see the same gain at every frequency in an allpass, while a comb filter has widely varying gains





(a)
$$H(z) = \frac{0.7 + z^{-7}}{1 + 0.7z^{-7}}$$
 (b) $H(z) = \frac{1}{1 + 0.7z^{-7}}$

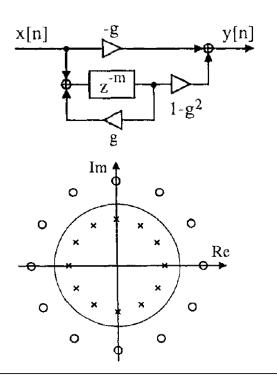


- The decay between successive samples in comb and allpass filters is described by the gain coefficient g_i
- In order for the comb filter's decay to correspond to a given reverberation time, we must have

$$\frac{20\log_{10}(g_i)}{m_i T} = \frac{-60}{T_r}$$

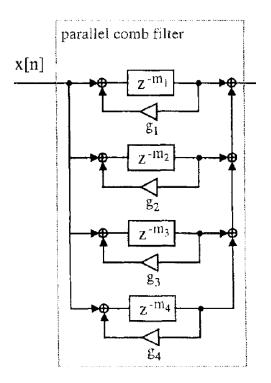
Combination of comb filters

- Single comb filters do not provide sufficient echo density
- In order to improve the echo density, we need to combine multiple comb filters
 - Cascading comb filters corresponds to multiplying their transfer functions
 - Frequency peaks not shared by all comb filters are cancelled by multiplication



Combination of comb filters

- Better to place comb filters in parallel
 - □ Example





Parallel comb filters

■ The poles of comb filters are given by

$$\prod_{i=1}^{N} (g_i - z^{m_i}) = 0$$

■ The poles have the same magnitudes

$$\gamma_i = \sqrt[m]{g_i} = 10^{-3T/T_r}$$

■ The modal density (No. of modes per Hz) is

$$D_m = \sum_{i=0}^{N-1} \tau_i = N \cdot \tau$$

100

Parallel comb filters

- Modal density turns out to be the same at all frequencies, unlike real rooms
- Above a threshold frequency, the modal density is constant
- The modal density of the comb filters is then set to the modal density above the threshold frequency

$$\sum_{i} \tau_{i} = D_{m} > D_{f} \approx \frac{T_{\text{max}}}{4}$$



Parallel comb filters

The echo density of the comb filters is approximatively given by

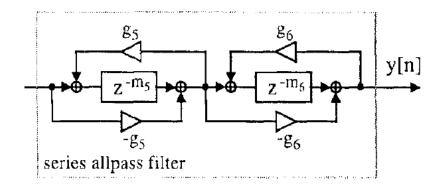
$$D_e = \sum_{i=0}^{N-1} \frac{1}{\tau_i} \approx \frac{N}{\tau}$$

Relating echo density and modal density provides:

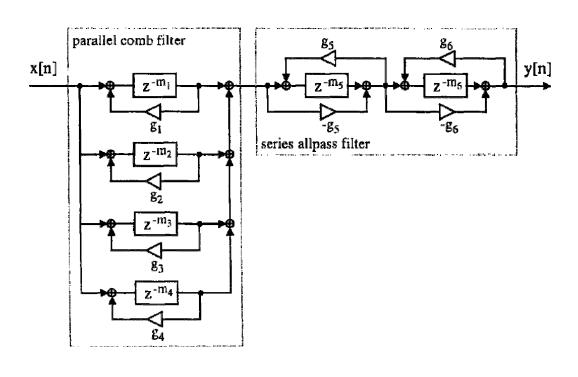
$$N \approx \sqrt{D_m D_e}$$

Combination of allpass filters

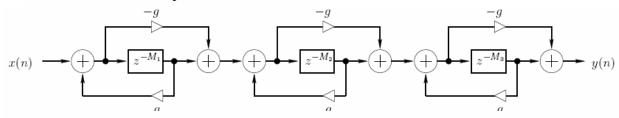
- Unlike comb filters, allpass filters must be cascaded
 - ☐ Multiplying freq. responses corresponds to adding phase responses



Schroeder's reverberator (1)



Schroeder Allpass Sections



- Typically, g = 0.7
- Delay-line lengths M_i mutually prime, and span successive orders of magnitude e.g., 1051, 337, 113
- Allpass filters in series are allpass
- Each allpass expands each nonzero input sample from the previous stage into an entire infinite allpass impulse response
- Allpass sections may be called "impulse expanders", "impulse diffusers" or simply "diffusers"
- NOT a physical model of diffuse reflection, but single reflections are expanded into many reflections, which is qualitatively what is desired.



Schroeder's reverberator

- Delays of the comb and allpass filters are chosen so that the ratio of the largest and smallest delay is 1.5 (typically 30 and 45 ms)
- The gains g_i of the comb filters are chosen to provide a desired reverberation time T_r according to

$$g_i = 10^{-3m_i T/T_v}$$

Allpass filters delays are set to 5 and 1.7 ms

A Schroeder Reverberator called JCRev RevIn $AP_{1051}^{0.7}$ $AP_{337}^{0.7}$ $AP_{113}^{0.7}$ $AP_{113}^{0.7}$ $FFCF_{5399}^{0.715}$ $FFCF_{5801}^{0.697}$ $FFCF_{5801}^{0.697}$ C

A Schroeder Reverberator called JCRev

• Three Schroeder allpass sections:

$$\mathsf{AP}_N^g \stackrel{\Delta}{=} \frac{g + z^{-N}}{1 + gz^{-N}}$$

• Four feedforward comb-filters:

$$\mathsf{FFCF}_N^g \stackrel{\Delta}{=} g + z^{-N}$$

• Schroeder suggests a progression of delays close to

$$M_i T \approx \frac{100 \text{ ms}}{3^i}, \quad i = 0, 1, 2, 3, 4.$$

- Comb filters impart distinctive coloration:
 - Early reflections
 - Room size
 - Could be one tapped delay line

A Schroeder Reverberator called JCRev

- Usage: Instrument adds scaled output to RevIn
- Reverberator output RevOut goes to four delay lines
 - Four channels decorrelated
 - Imaging of reverberation between speakers avoided
- For stereo listening, Schroeder suggests a mixing matrix at the reverberator output, replacing the decorrelating delay lines
- A mixing matrix should produce maximally rich yet uncorrelated output signals

Feedback Delay Networks... ...at a glance

- Unitary matrix: definition
 - □ A matrix is unitary if :

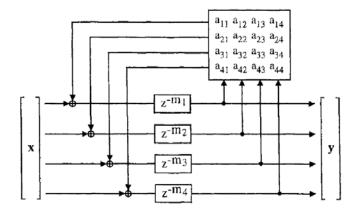
$$\|\mathbf{M} \cdot \mathbf{u}\| = \|\mathbf{u}\|$$

□ We can also write that a matrix is unitary if

$$||\mathbf{M}\mathbf{M}^T|| = ||\mathbf{M}^T\mathbf{M}|| = 1$$



FDN



$$\mathbf{A} = \begin{bmatrix} a_{11} & a_{12} & a_{13} & a_{14} \\ a_{21} & a_{22} & a_{23} & a_{24} \\ a_{31} & a_{32} & a_{33} & a_{34} \\ a_{41} & a_{42} & a_{43} & a_{44} \end{bmatrix}$$

- \square Stability of the feedback loop is guaranteed if A = gM where M is an unitary matrix and |g|<1.
- □Outputs will be mutually incoherent: we can use the FDN to render the diffuse soundfield with a 4 loudspeaker system.
- ☐ The early reverbeartions can be simulated by appropriately injecting the input signal into the delay lines.



Jot's reverberator

$$\mathbf{A} = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1N} \\ a_{21} & a_{22} & \cdots & a_{2N} \\ \vdots & \vdots & \vdots & \vdots \\ a_{N1} & a_{42} & \cdots & a_{4N} \end{bmatrix} \quad \mathbf{b} = \begin{bmatrix} b_1 \\ \vdots \\ b_N \end{bmatrix} \quad \mathbf{c} = \begin{bmatrix} c_1 \\ \vdots \\ c_N \end{bmatrix}$$



Jot's reverberator

The input-output relation of Jot's reverberator is given by

$$y(z) = \mathbf{c}^T \mathbf{s}(z) + dx(z)$$

$$\mathbf{s}(z) = \mathbf{D}(z)[\mathbf{A}\mathbf{s}(z) + \mathbf{b}x(z)]$$

with
$$\mathbf{s}(z) = \begin{bmatrix} s_1(z) \\ \vdots \\ s_N(z) \end{bmatrix}$$
 and $\mathbf{D}(z) = \begin{bmatrix} z^{-m_1} & 0 \\ & \ddots & \\ 0 & z^{-m_N} \end{bmatrix}$



Jot's reverberator

System transfer function:

$$H(z) = \frac{y(z)}{x(z)} = \mathbf{c}^T [\mathbf{D}(z^{-1}) - \mathbf{A}]^{-1} \mathbf{b} + d$$

Zeros:

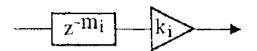
$$\det[\mathbf{A} - \frac{\mathbf{b}\mathbf{c}^T}{d} - \mathbf{D}(z^{-1})] = 0$$

■ Poles:

$$\det[\mathbf{A} - \mathbf{D}(z^{-1})] = 0$$

Jot's reverberator

- Moorer noted that convolving exponentially decaying white noise with source signals produces a very natural sounding.
- As a consequence, by introducing absorptive losses into a lossless prototype, we should obtain a natural sounding reverberator.
- This is accomplished by associating a gain with each delay:





Jot's reverberator

The logarithm of the gain is proportional to the length of the delay:

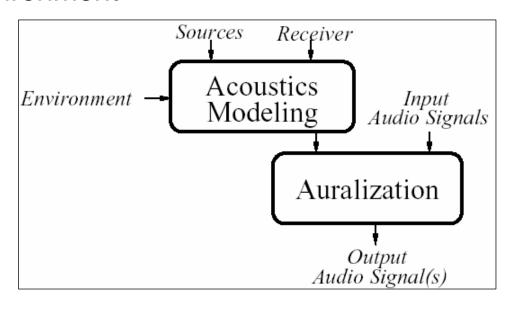
$$k_i = \gamma^{m_i}$$

- The above modification has the effect of replacing z with z/γ in the transfer function
- The lossless prototype response h[n] will be multiplied by an exponential envelope γ^n

Modeling the Environment

Modeling the environment

 Simulate reverberations due to environment





Motivations

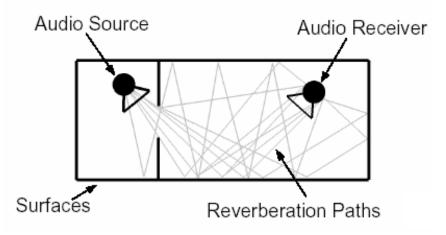
Acoustical environment provides ...

- Sense of presence
- Comprehension of space
- Localization of auditory cues
- Selectivity of audio signals ("cocktail party effect")



Geometric acoustic modeling

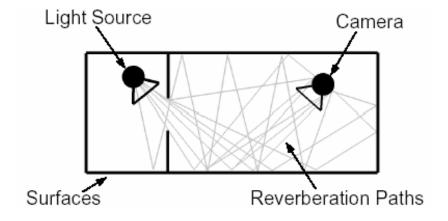
Spatialize sound by computing reverberation paths from source to receiver





Similarities to Graphics

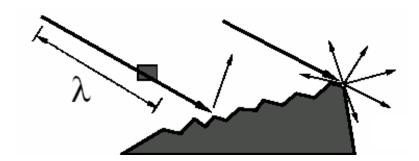
Both model wave propagatation





Differences from Graphics I

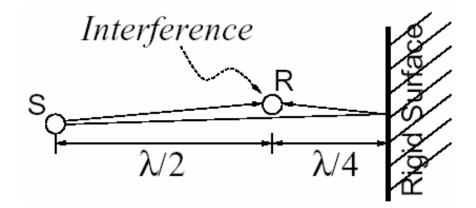
- Sound has longer wavelengths than light
 - □ Diffractions are significant
 - □ Specular reflections dominate diffuse reflections
 - □ Occlusions by small objects have little effect





Differences from Graphics II

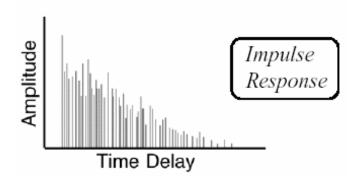
- Sound waves are coherent
 - ☐ Modeling phase is important





Differences from Graphics III

- Sound travels more slowly than light
 - □ Reverberations are perceived over time





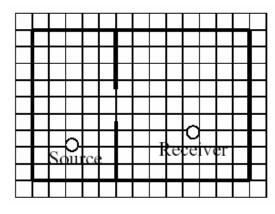
Overview of approaches

- Finite element methods
- Boundary element methods
- Image source methods
- Ray tracing
- Beam tracing



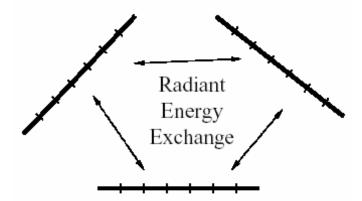
Finite element methods

Solve wave equation over grid-aligned mesh



Boundary element methods

Solve wave equation over discretized surfaces





Boundary Element Trade-offs

- Advantages
 - Works well for low frequencies
 - □ Simple formulation

$$B_i = E_i + \rho_i \sum B_j F_{ij}$$

Boundary Element Trade-offs

Disadvantages

- □ Complex function stored with each element
- □ Form factors must model diffractions & specularities
- ☐ Elements must be much smaller than wavelength

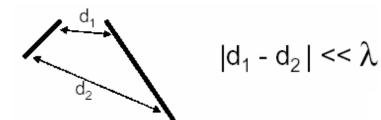


Image source methods

Consider direct paths from "virtual sources"

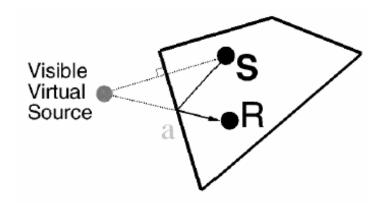




Image source trade-offs

Advantages

☐ Simple for rectangular rooms

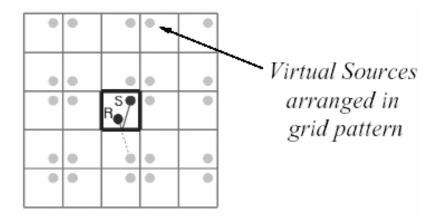


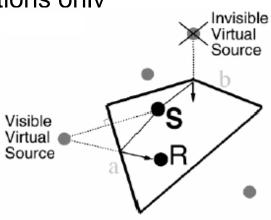


Image source trade-offs

Disadvantages

 \Box O(n^r) visibility checks in arbitrary environments

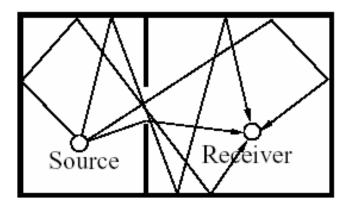
☐ Specular reflections only





Path tracing

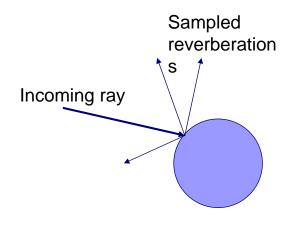
Trace paths between source and receiver





Path Tracing Trade-offs

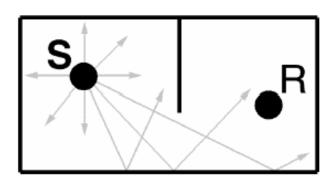
- Advantages
 - ☐ Models all types of surfaces and scattering
 - ☐ Simple to implement





Path Tracing Disadvantages

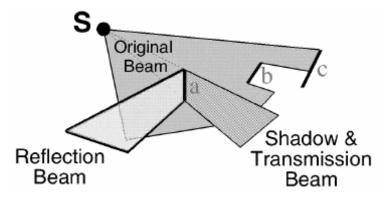
- Disadvantages
 - ☐ Subject to sampling errors (aliasing)
 - □ Depends on receiver position





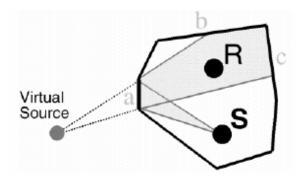
Beam Tracing

Trace beams (bundles of rays) from source



Beam Tracing Trade-offs

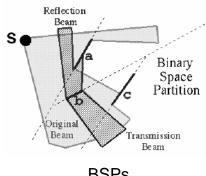
- Advantages
 - □ Takes advantage of spatial coherence
 - □ Predetermines visible virtual sources



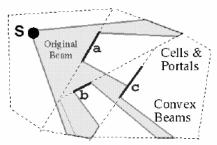


Beam Tracing Disadvantages

- Disadvantages
 - □ Difficult for curved surfaces or refractions
 - □ Requires efficient polygon sorting and intersection



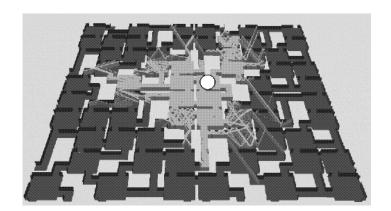
BSPs



Cell adjacency graphs

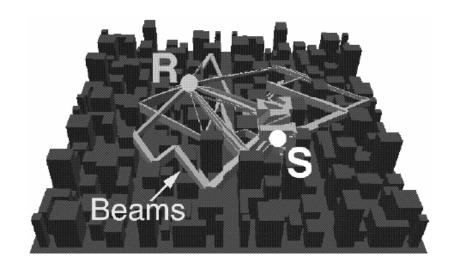
Complex 3D Environments

Precompute beam tree for stationary source



Interactive Performance

Lookup beams containing moving receiver





Summary

- FEM/BEM
 - □ best for low frequencies
- Image source methods
 - □ best for rectangular rooms (very common)
- Path tracing
 - □ best for high-order reflections (very common)
- Beam tracing
 - □ best for precomputation